# Doing Quantum Physics with Split-Complex Numbers

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Mathematics-Physics Joint Colloquium, February 21, 2014

### Outline



2 Complex Numbers are for Light

### 3 Quantum Physics with Split-Complex Numbers

# Path Integral

 In quantum mechanics, the amplitude of a particle to propagate from a point q<sub>I</sub> to a point q<sub>F</sub> in time T is given by

$$\langle q_F | e^{-\frac{i}{\hbar}\hat{H}T} | q_I \rangle = \int Dq(t) e^{\frac{i}{\hbar}\int_0^T dt L(\dot{q},q)}$$

•  $L(\dot{q},q)$  is the Lagrangian

$$L(\dot{q},q) = \frac{m}{2}\dot{q}^2 - V(q)$$

• Dq(t) is the Feynman measure given by

$$\int Dq(t) := \lim_{N \to \infty} \left( \frac{-im\hbar}{2\pi\delta t} \right)^{\frac{N}{2}} \left( \prod_{k=1}^{N-1} \int dq_k \right)$$

where  $\delta t = \frac{T}{N}$ .

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### Euclideanisation

- Wick rotation t → it turns Minkowski spacetime into Euclidean spacetime.
- Accordnigly, the path integral turns into Euclidean path integral

$$\langle q_F | e^{-rac{i}{\hbar}\hat{H}T} | q_I 
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# Problems with Euclideanisation

- It is troublesome that path integral cannot be calculated in actual spacetime and that it must be calculated in Euclidean spacetime which is not physical spacetime.
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# Example

• Define a function 
$$f:(-\infty,0)\longrightarrow (-\infty,\infty)$$
 by

$$f(x) = \sum_{n=1}^{\infty} e^{nx}$$
$$= e^{x} + e^{2x} + \cdots$$

• Since  $|e^x| < 1$  on  $(-\infty, 0)$ , f(x) converges to

$$f(x) = \frac{1}{e^{-x} - 1}$$

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#### Example Continued

• The Wick rotated  $(x \mapsto ix)$  function

$$g(x) = \sum_{n=1}^{\infty} e^{inx}$$
$$= \sum_{n=1}^{\infty} [\cos(nx) + i\sin(nx)]$$

does not converge.

• For instance,  $g(-2\pi) = \infty$ , while  $f(-2\pi) = \frac{1}{e^{2\pi}-1}$ .

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# Why Complex Numbers in Quantum Mechanics?

- Light must be described by electromagnetic waves or by particles (Wave-Particle Duality)
- de Broglie hypothesised that what is true for photons should be valid for any particle.
- A photon can be described by the complex plane wave

$$\psi(\mathbf{r},t) = A \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$

with energy E and momentum vector  $\mathbf{p}$  satisfying the equations

 $E = \hbar \omega, \ \mathbf{p} = \hbar \mathbf{k}$ 

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# Electric-Magnetic Duality

• Maxwell's equations in vacuum are:

$$abla \cdot \mathbf{B} = 0, \ \nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$
  
 $abla \cdot \mathbf{E} = 0, \ \nabla \times \mathbf{B} - \frac{\partial \mathbf{E}}{\partial t} = 0$ 

• The transformation

$$\mathsf{B}\mapsto\mathsf{E},\ \mathsf{E}\mapsto-\mathsf{B}$$

takes the first pair of equations to the second and vice versa. This symmetry is called *Electric-Magnetic Duality*.

• The duality hints that the electric and magnetic fields are part of a unified whole, the *electromagnetic field*.

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# Electromagnetic Field as a Complex-Valued Vector Field

#### • Let us introduce a complex-valued vector field

#### $\mathscr{E}=\mathbf{E}+i\mathbf{B}$

• The duality amounts to the transformation

 $\mathscr{E}\mapsto -i\mathscr{E}$ 

• The vacuum Maxwell's equations boil down to two equations for  $\mathscr{E}$ :

$$\nabla \cdot \mathscr{E} = 0, \ \nabla \times \mathscr{E} = i \frac{\partial \mathscr{E}}{\partial t}$$

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### Plane Wave as Electromagnetic Field

Let **k** be a vector in  $\mathbb{R}^3$  and let  $\boldsymbol{\omega} = |\mathbf{k}|$ . Fix  $\mathbf{E} \in \mathbb{C}^3$  with  $\mathbf{E} \cdot \mathbf{k} = 0$ and  $\mathbf{E} \times \mathbf{k} = i\boldsymbol{\omega}\mathbf{E}$ . Then the plane wave

$$\mathscr{E}(\mathbf{r},t) = \mathbf{E} \exp[i(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$

satisfies the vacuum Maxwell's equations.

# The Light Cone and Two-Sphere

$$\mathbb{N}^3 = \left\{ (t, x, y, z) \in \mathbb{R}^{3+1} : t^2 - x^2 - y^2 - z^2 = 0 \right\}$$

- Let  $\mathbb{N}^3_+$  and  $\mathbb{N}^3_-$  denote the future and the past light cones respectively. The multiplicative group  $\mathbb{R}^+$  acts on  $\mathbb{N}^3_+$  and  $\mathbb{N}^3_-$  respectively by scalar multiplication.
- Define  $f_{\pm}: \mathbb{N}^3_{\pm} \longrightarrow S^2$  by  $f_{\pm}(t, x, y, z) = \left(\frac{x}{t}, \frac{y}{t}, \frac{z}{t}\right)$ . Then  $f_{\pm}$  are continuous surjections i.e. identification maps.
- The orbit spaces N<sup>3</sup><sub>+</sub>/ℝ<sup>+</sup> and N<sup>3</sup><sub>-</sub>/ℝ<sup>+</sup> are identified with the two-sphere S<sup>2</sup>. The identification is a homeomorphism. It is indeed a diffeomorphism.

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Celestial Sphere and Complex Numbers

- For an observer at the origin (the event), light rays through his eye correspond to null lines through the origin.
- The past null directions constitute the field of vision of the observer which is the two-sphere S<sup>2</sup>.
- The two-sphere S<sup>2</sup> is the extended complex plane ℂ∪{∞} called the *Riemann sphere*.

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#### In the beginning, God might have said

"Let there be complex numbers!"

#### Wave Functions are real?

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- What if we assume that wave functions are real, say they represent actual waves in spacetime?

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Split-Complex Number System

• Let  $\mathbb{C}'$  be a real commutative algebra spanned by 1 and *j*, with multiplication law:

$$1 \cdot j = j \cdot 1 = j, \ j^2 = 1$$

An element of  $\mathbb{C}' = 1\mathbb{R} \oplus j\mathbb{R}$  is called a *split-complex number*, a *paracomplex number*, or a *hyperbolic number*.

•  $\zeta \in \mathbb{C}'$  is uniquely expressed as  $\zeta = x + jy$ . The conjugate  $\overline{\zeta}$  is defined by  $\overline{\zeta} = x - jy$  and the squared modulus  $|\zeta|^2$  is defined to be

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•  $\mathbb{C}'$  is identified with  $\mathbb{R}^{1+1}$ .

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## Algebraic Representation of $\mathbb{R}^{3+1}$

• The spacetime  $\mathbb{R}^{3+1}$  can be identified with a set of  $2\times 2$  Hermitian matrices:

$$X = (t, x, y, z) \longleftrightarrow \underline{X} = \begin{pmatrix} t + jz & x + iy \\ x - iy & t - jz \end{pmatrix} = te_0 + xe_1 + ye_2 + jze_3$$

where

$$\mathbf{e}_0 = \left(\begin{array}{cc} 1 & 0 \\ 0 & 1 \end{array}\right), \ \mathbf{e}_1 = \left(\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array}\right), \ \mathbf{e}_2 = \left(\begin{array}{cc} 0 & i \\ -i & 0 \end{array}\right), \ \mathbf{e}_3 = \left(\begin{array}{cc} 1 & 0 \\ 0 & -1 \end{array}\right)$$

are Pauli spin matrices.

• The identification is an isometry:

$$\langle X, Y \rangle = \frac{1}{2} \operatorname{tr}(\underline{XY}^{\dagger})$$

In particular,  $|X|^2 = \det \underline{X}$ .

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$$\longleftrightarrow (t + jz) + (x + iy) \in \mathbb{C}' \oplus \mathbb{C}$$

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$$\mathbb{R}^{3+1} \cong \mathbb{C}' \oplus \mathbb{C}$$

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#### Euler's Formula

 $\bullet\,$  In  $\mathbb{C}',$  there is an analogue of the Euler's formula:

$$\exp(j heta) = \cosh heta + j \sinh heta$$

where  $-\infty < \theta < \infty$ . The number  $\theta$  is called a *hyperbolic* angle.

- $\exp(j\theta)$  is a point on the hyperbola  $x^2 y^2 = 1$ .
- In matrix form, exp(jθ) can be written as

$$\left(\begin{array}{c}\cosh\theta&\sinh\theta\\\sinh\theta&\cosh\theta\end{array}\right)\in\mathrm{SO}^+(1,1)$$

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#### Split-Complex Plane Wave

- Let us consider a split-complex plane wave  $\psi(\mathbf{r}, t) = A \exp[j(\mathbf{k} \cdot \mathbf{r} - \omega t)]$ , where A is a real number.
- If we assume that the wave is traveling at the speed of light in vacuum,  $\psi({\bf r},t)$  satisfies the wave equation

$$-\frac{1}{c^2}\frac{\partial^2\psi}{\partial t^2} + \nabla^2\psi = 0$$

• The energy operator  $\hat{E}$  and the momentum operator  $\hat{p}$  are obtained as

$$\hat{E} = -j\hbar \frac{\partial}{\partial t}, \ \hat{p} = j\hbar \nabla$$

$$-j\hbar\frac{\partial\psi}{\partial t} = \frac{\hbar^2}{2m}\nabla^2\psi$$

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- If we assume that the wave is traveling at the speed of light in vacuum,  $\psi(\mathbf{r},t)$  satisfies the wave equation

$$-\frac{1}{c^2}\frac{\partial^2\psi}{\partial t^2}+\nabla^2\psi=0$$

• The energy operator  $\hat{E}$  and the momentum operator  $\hat{p}$  are obtained as

$$\hat{E} = -j\hbar \frac{\partial}{\partial t}, \ \hat{p} = j\hbar \nabla$$

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#### Negative Probability?

• Let 
$$\psi^+(\mathbf{r},t) = A \exp[j(\mathbf{k} \cdot \mathbf{r} - \omega t)]$$
 and  
 $\psi^-(\mathbf{r},t) = Aj \exp[j(\mathbf{k} \cdot \mathbf{r} - \omega t)].$ 

•  $\psi^{-}(\mathbf{r}.t)$  also satisfies the Schrödinger equation.

• While 
$$|\psi^+(\mathbf{r},t)|^2 = A^2 > 0$$
,  $|\psi^-(\mathbf{r},t)|^2 = -A^2 < 0$ .

- The negative sign may be interpreted as a difference in sign of unit charge between a particle and its antiparticle.
- If  $\psi^+(\mathbf{r},t)$  is a plane wave for a particle with charge density  $\rho_e^+ = e\overline{\psi^+}\psi^+$ , then  $\psi^-(\mathbf{r},t)$  may be considered as a plane wave for its antiparticle with charge density  $\rho_e^- = e\overline{\psi^-}\psi^-$ .

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### $\psi^+(\mathbf{r},t)$ and $\psi^-(\mathbf{r},t)$

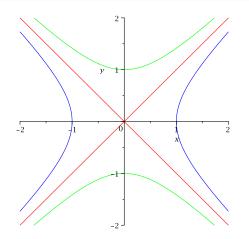


Figure :  $\psi^+$  (in blue) and  $\psi^-$  (in green)

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#### Split-Complex Structure and the Charge Conjugation Map

• Define a linear endomorphism  $\mathscr{J}:\mathbb{C}'\longrightarrow\mathbb{C}'$  by

$$\mathcal{J}1=j, \mathcal{J}j=1$$

•  $\mathscr{J}$  satisfies

$$\mathcal{J}^2 = \mathcal{I}, \ \langle \mathcal{J} \zeta_1, \mathcal{J} \zeta_2 \rangle = - \langle \zeta_1, \zeta_2 \rangle$$

Thus  $\mathscr{J}$  is an anti-isometry.  $\mathscr{J}$  is called the *associated split-complex structure* of  $\mathbb{C}'$ .

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Two Hilbert Spaces  $\mathscr{H}^+$  and  $\mathscr{H}^-$ 

- {ψ<sub>n</sub><sup>+</sup>(**r**, t) : ψ<sub>n</sub><sup>+</sup>(**r**, t) = A<sub>n</sub> exp[j(**k**<sub>n</sub> · **r** − ω<sub>n</sub>t)], n = 1, 2, · · · } forms a countable basis for a spilt-complex Hilbert space ℋ<sup>+</sup> over real field ℝ.
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#### Twin Universes

- Under the interpretation, it appears that antiparticles are living in a different spacetime,  $\mathbb{R}^{3+1}(t, x, y, z)$  with metric signature (-+--).
- Big Bang might have created twin (not identical though) universes, one made of matter and the other made of antimatter.
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#### Path Integral Redux

 The amplitude of a particle to propagate from a point q<sub>I</sub> to a point q<sub>F</sub> in time T is obtained as

$$\langle q_{\mathsf{F}}|e^{-rac{j}{\hbar}\hat{H}\mathcal{T}}|q_{\mathsf{I}}
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• The Feynman meassure Dq(t) is given by

$$\int Dq(t) := \lim_{N \to \infty} \left( \frac{2\pi m \hbar j}{\delta t} \right)^{\frac{N}{2}} \left( \prod_{k=1}^{N-1} \int dq_k \right)$$

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